# Behaviour near extinction for the fast diffusion equation in bounded domains 

Matteo Bonforte

Departamento de Matemáticas,
Universidad Autónoma de Madrid,
Campus de Cantoblanco
28049 Madrid, Spain
email: matteo.bonforte@uam.es
http://www.uam.es/matteo.bonforte
(Joint work with G. Grillo and J. L. Vázquez)

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## The Dirichlet Problem for the Fast Diffusion Equation in $\Omega \subset \mathbb{R}^{d}$

We consider, in a bounded and smooth domain $\Omega$, positive solutions to:

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\begin{cases}\partial_{\tau} u=\Delta\left(u^{m}\right)=\nabla \cdot\left(u^{m-1} \nabla u\right), & \forall(\tau, y) \in(0,+\infty) \times \Omega \\ u(0, y)=u_{0}, & \forall y \in \Omega \\ u(\tau, y)=0, & \forall(\tau, y) \in(0,+\infty) \times \partial \Omega\end{cases}
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\text { where } 0<m<1 \text { (i.e. Fast Diffusion, FDE ) }
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- Existence and uniqueness of weak solutions for the parabolic problem is well known for any $m>0$. Recall that $0<m<1$ is the Fast Diffusion case, $m=1$ is the Linear Heat Equation and $m>1$ is the Porous Medium case.
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0 \leq u_{0} \in \mathrm{~L}^{r}(\Omega) \quad \text { with } \quad r \geq 1 \quad \text { and } \quad r>\frac{d(1-m)}{2}
$$

so that the corresponding solution is bounded and nonnegative for all $m>0$.

## Some Properties of Solutions

- Since we deal with the Fast Diffusion case $m<1$, the mass $\int_{\Omega} u(y, \tau) \mathrm{d} y$ is not preserved, and solutions extinguish in finite time

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\exists T=T\left(u_{0}\right): u(\tau, \cdot) \equiv 0 \quad \forall t \geq T
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Consequence of Sobolev and Poincaré inequalities (sufficient condition).

- Under our hypothesis, solutions are indeed positive in $\Omega \times(0, T)$ and for all $0<m<1$, as a consequence of parabolic (intrinsic) Harnack inequalities:
- For $\frac{d-2}{d}<m<1$. DiBenedetto et al. (1992)
- For all $0<m<1$, Bonforte and Vázquez (2010) and they are at least $C^{\alpha}(\bar{\Omega})$ (DiBenedetto et al. 1988, 1992).
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## Review of previous results

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The properties of the rescaled problem are related to the stationary equation


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The crucial exponent is

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m_{s}=\frac{d-2}{d+2} ; \quad \text { we shall consider the range } \quad m_{s}<m<1
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## (First Pioneering Result) J. G. Berryman, C. J. Holland ARMA (1980)

Let $m_{s}<m<1$. Then there exists a sequence of times $t_{n} \rightarrow \infty$ as $n \rightarrow \infty$ and one or several solutions $S$ to the stationary problem such that

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v\left(t_{n}\right) \xrightarrow[n \rightarrow \infty]{W_{0}^{1,2}(\Omega)} S
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(Uniqueness of asymptotic profile) E. Feiresl, F. Simondon J. Dynamic Diff. Eq. (2000) Let $v, S$ be as above and assume $m_{s}<m<1$. Then there exists a unique stationary solution $S$ such that

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## (Global Harnack Principle)

E. DiBenedetto, Y. C. Kwong, V. Vespri Indiana Univ. Math. J. (1991)

Let $w$ be the solution to the rescaled Dirichlet problem with $m_{s}<m<1$. Then, for any $\sigma>0$ there exist positive constants $\lambda, \mu>0$ depending on $d, m,\left\|u_{0}\right\|_{m+1},\left\|\nabla u_{0}^{m}\right\|_{2}, \partial \Omega$ and $\sigma$, such that for any $t \geq \sigma$ and for any $x \in \Omega$

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\lambda \operatorname{dist}(x, \partial \Omega)^{1 / m} \leq v(t, x) \leq \mu \operatorname{dist}(x, \partial \Omega)^{1 / m}
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The constants $\lambda, \mu$ may deteriorate when $m \rightarrow 1$ or $m \rightarrow m_{s}$.
(Convergence in Relative Error) M.B., G. Grillo, J.L. Vázquez, JMPA (2011)
Let $u$ be the solution to the Dirichlet problem and $T=T\left(m, d, u_{0}\right)$ be its extinction time. Then we have that

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\lim _{\tau \rightarrow T^{-}}\left\|\frac{u(\tau, \cdot)}{\mathcal{U}(\tau, \cdot)}-1\right\|_{L^{\infty}(\Omega)}=0
$$

where the special solution $\mathcal{U}$ is defined as

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\mathcal{U}(\tau, x)=S(x)[(T-\tau) / T]^{1 /(1-m)} \quad\left[\text { one has } S(x) \sim \operatorname{dist}(x, \partial \Omega)^{\frac{1}{m}}\right]
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and $S$ is a suitable positive classical solution to the stationary problem. Equivalently, the following improved Global Harnack Principle


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Steps of the proof.

- Consider the function $\phi=\frac{\nu^{m}}{S^{m}}-1$. Then it satisfies the equation

where $F$ is given by $F(\phi)=\mathbf{c}\left[(1+\phi)^{1 / m}-(1+\phi)\right]$.
- Convergence far away from the boundary is easy.
- One can choose positive constants $A, B, C$ and $t_{0}$, so that the function

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\Phi(t, x)=C-B d(x)-A\left(t-t_{0}\right)
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is a supersolution to the differential equation satisfied by $\phi$, in a small neighborhood of the spatial boundary $\Omega_{\delta}=:\{x \in \Omega: \operatorname{dist}(x, \partial \Omega)<\delta\}$. Technical.

- Use parabolic comparison to compare $\phi$ and $\Phi$ in $t \in\left(t_{0}, T\right] \times \Omega_{\delta}$. $\square$

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## The fast diffusion equation on bounded domains

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In the sequel, the constants $m_{\sharp}, \gamma_{0}$ are explicit. They depend on $m$ and on the geometry of the domain.
(Decay Rates, Rescaled Vession) M.B., G.Grillo, J.L. Maguen, JMPA (20)I)
Let $m_{\sharp}<m<1$. Let $v$ be the rescaled solution corresponding to an initial datum $u_{0}$, and let $S$ be the stationary profile to which the solution converges. Let $0<\gamma<\gamma_{0}$. Then for all $t>t_{0}$ :
where $\bar{\theta}(t)$ is the mean of $\theta(t)$ w.r.t. to the measure $S^{m+1} \mathrm{~d} x$.
Therefore the following holds:

$$
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Finally, for all $q \in(0, \infty]$ :

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In the sequel, the constants $m_{\sharp}, \gamma_{0}$ are explicit. They depend on $m$ and on the geometry of the domain.

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Let $m_{\sharp}<m<1$. Let $v$ be the rescaled solution corresponding to an initial datum $u_{0}$, and let $S$ be the stationary profile to which the solution converges. Let $0<\gamma<\gamma_{0}$. Then for all $t>t_{0}$ :

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\mathcal{E}[\theta(t)]:=\frac{1}{2} \int_{\Omega}|\theta(t)-\bar{\theta}(t)|^{2} S^{m+1} \mathrm{~d} x \leq \mathrm{e}^{-\gamma\left(t-t_{0}\right)} \mathcal{E}\left[\theta\left(t_{0}\right)\right]
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## The fast diffusion equation on bounded domains

## 0000000000000000000000

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## (Decay Rates, Original Variables)

Let $\max \left\{m_{\sharp}, m_{c}\right\}<m<1$. Let $u$ be the solution to the original FDE Problem, let $T=T\left(m, d, u_{0}\right)$ be its extinction time, and let $\mathcal{U}_{T}$ be previous special solution, so that $u(\tau) / \mathcal{U}_{T}(\tau) \rightarrow 1$ uniformly as $\tau \rightarrow T$. Then, for any $\bar{\gamma}<\bar{\gamma}_{0}:=\gamma_{0} T$ there exists a constant $\kappa_{0}>0$ such that

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The weighted convergence of (2) is somehow stronger than the non-weighted $\mathrm{L}^{p}$ - norm convergence, since the weight $S^{m-1}$ is singular at the boundary.

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Similar consideration also work for the porous media case ( $m>1$ ), which has been studied long ago by completely different methods (Aronson-Peletier, JDE (1981)).

## (Decay Rates, Porous Medium)

Let $m>1$, let $v$ be a the rescaled solution, that converges to its unique stationary state $S$, and let $\theta=v / S$. Then, for all $0<\beta<2+\frac{K m}{m-1}$ there exists a time $t_{1}$ depending on $m, d, \beta$ and on the constant $K>0$ of the weighted Poincaré inequality, such that the entropy decays as

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where the special solution $\mathcal{U}$ is defined by $\mathcal{U}(\tau, x)=S(x)(1+\tau)^{-1 /(m-1)}$.

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\mathcal{E}[\theta(t)] \leq \mathcal{E}\left[\theta\left(t_{1}\right)\right] \mathrm{e}^{-\beta\left(t-t_{1}\right)} \quad \text { for all } t \geq t_{1} \tag{3}
\end{equation*}
$$

Moreover for all $q \in(0, \infty]$

$$
\|v(t, \cdot)-S(\cdot)\|_{\mathrm{L}^{q}(\Omega)} \leq \kappa_{1} \mathrm{e}^{-\left(t-t_{1}\right)} \quad \text { for all } t \geq t_{1}
$$

In original variables we obtain that for all $q \in(0, \infty]$

$$
\|u(\tau, \cdot)-\mathcal{U}(\tau, \cdot)\|_{\mathrm{L}^{q}(\Omega)} \leq \frac{\kappa_{2}}{(1+\tau)^{1+\frac{1}{m-1}}}
$$

where the special solution $\mathcal{U}$ is defined by $\mathcal{U}(\tau, x)=S(x)(1+\tau)^{-1 /(m-1)}$.

## Short review on the linear case.

Consider the homogeneous Dirichlet problem for the linear heat equation $u_{\tau}=\Delta u$.

- Rescale $v(x, t)=e^{\lambda_{1} t} u(x, t)$ to get the equation $v_{t}=\Delta v+\lambda_{1} v$.
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- The equation for the relative error $\theta=v / \Phi_{1}-1$ is $\theta_{t}=\Phi_{1}^{-2} \nabla \cdot\left(\Phi_{1}^{2} \nabla \theta\right)$
- The so-called Dirichlet Laplacian has purely discrete spectrum. Let $\lambda_{j}, \Phi_{j}, j=$
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The spectral representation for the heat semigroup gives

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u(x, t)=\sum_{j=1}^{\infty} c_{j} e^{-\lambda_{j} t} \Phi_{j}(x) \quad \text { with } c_{j}=\int_{\Omega} u_{0} \Phi_{j} \mathrm{~d} x
$$ so that



In other words, the solution $u(t)$ behaves like $U_{1}(x, t)=c_{1} \mathrm{e}^{-\lambda_{1} t} \Phi_{1}$ and the relative error $\theta$, decays exponentially in time with a rate $\lambda_{2}-\lambda_{1}$. (recall that $\Phi_{2} / \Phi_{1}$ is bounded.)

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\theta:=\frac{u}{c_{1} \mathrm{e}^{-\lambda_{1} t} \Phi_{1}}-1 \underset{t \rightarrow+\infty}{\sim} \frac{c_{2}}{c_{1}} \frac{\Phi_{2}}{\Phi_{1}} \mathrm{e}^{-\left(\lambda_{2}-\lambda_{1}\right) t}
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## The fast diffusion equation on bounded domains

## 000000000000000000000

Short review on the linear case

## Short review on the linear case (continued).

- In the nonlinear setting, no spectral representation is available. It is natural to investigate the behaviour of $\theta$ by working in the weighted space $L^{2}\left(\Phi_{1}^{2} \mathrm{~d} x\right)$, where the weighted mean is preserved:

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\frac{\mathrm{d}}{\mathrm{~d} t} \int_{\Omega} \theta \Phi_{1}^{2} \mathrm{~d} x=\int_{\Omega} \nabla \cdot\left(\Phi_{1}^{2} \nabla \theta\right) \mathrm{d} x=0 .
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Then we notice that:


We shall assume that $\theta_{\Phi_{1}}=0$, where $g_{\Phi_{1}}=\left(\int_{\Omega} g \Phi_{1}^{2} \mathrm{~d} x\right) /\left(\int_{\Omega} \Phi_{1}^{2} \mathrm{~d} x\right)$

- In order to get a decay rate for $E[\theta]=\int_{\Omega} \theta^{2} \Phi_{1}^{2} \mathrm{~d} x$ we need the following intrinsic Poincaré inequality: for all $f \in W_{0}^{1,2}(\Omega)$ and $g=f / \Phi_{1}$, we have

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$$
\frac{\pi^{2}}{\operatorname{diam}(\Omega)^{2}}<\lambda_{2}-\lambda_{1} \leq \frac{d \pi^{2}}{\operatorname{inr}(\Omega)^{2}}
$$

This bounds can be improved when further geometrical properties of $\Omega$ hold.

Sketch of the proof.
Step 1: an "entropy functional" and its derivative. Recall that

$$
\mathcal{E}[\theta(t)]=\frac{1}{2} \int_{\Omega}|\theta(t)-\bar{\theta}(t)|^{2} S^{1+m} \mathrm{~d} x
$$

where $S$ is a (positive) solution to the elliptic problem

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\begin{cases}-\Delta S^{m}=\mathbf{c} S & \text { in } \Omega \\ S=0 & \text { on } \partial \Omega\end{cases}
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whenever $m_{s}<m<1$. We then have:

## (Entropy/Entropy-production)

Let $m_{s}<m<1$ and $\theta$ be the solution to the equation
$\theta_{t}=\frac{1}{S^{m+1}} \nabla \cdot\left(S^{2 m} \nabla(1+\theta)^{m}\right)+\mathbf{c} f(\theta), \quad$ with $\quad f(\theta)=(1+\theta)-(1+\theta)^{m}$.
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for all sufficiently large times, where $\varepsilon(t):=\|\theta(t, \cdot)\|_{\infty} \rightarrow 0$

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Step 2: a weighted Poincaré inequality. We prove the following:

## Poincaré inequalities

Let $f \in W_{0}^{1,2}(\Omega), \phi_{1}$ the ground state eigenfunction of the Dirichlet Laplacian, $g=$ $f / \phi_{1}$ and $S$ as above. Then the following inequality holds

where $\Lambda=\lambda_{2}-\lambda_{1}>0$ is the optimal constant in the intrinsic Poincaré inequality


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Step 2: a weighted Poincaré inequality. We prove the following:

## Poincaré inequalities

Let $f \in W_{0}^{1,2}(\Omega), \phi_{1}$ the ground state eigenfunction of the Dirichlet Laplacian, $g=$ $f / \phi_{1}$ and $S$ as above. Then the following inequality holds

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Hence it is necessary to get information on the ratio $\frac{k_{0}(m)^{2}}{k_{1}(m)^{2}}$ in order to get exponential decay for $\mathcal{E}$ from the above inequalities, at least when $m$ is close to one. Recall that $k_{0}, k_{1}$ are such that
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The difficult issue is to estimate the ratio $\frac{k_{0}(m)^{2}}{k_{1}(m)^{2}}$.
On the one hand, one can prove results about quantitative elliptic Harnack inequalities for the equation $-\Delta u=u^{p}$. This is the topic of M.B., G. Grillo, J.L. Vázquez (2011, in preparation).

The resulting bounds give explicit constants in the Harnack inequality. It is then possible to use them to compare solutions with different values of $p$, which then yield the required bounds on $\frac{k_{0}(m)^{2}}{k_{1}(m)^{2}}$.

Such bounds then yield explicit $m_{\sharp}$ and $\gamma_{0}$, but it has to be remarked that unfortunately

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\lim _{m \uparrow 1} \frac{k_{0}(m)^{2}}{k_{1}(m)^{2}}<1
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which is not what is expected.
But we can also prove the following purely elliptic result (see also Grossi (Annali Pisa, 2009) for related results):

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## Nonlinear elliptic problems near $p=1$

Let $p=1 / m$. Let $U_{p}$ be a family of solutions of the problem

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\begin{cases}-\Delta U=\lambda_{p} U^{p} & \text { in } \Omega  \tag{4}\\ U>0 & \text { in } \Omega \\ U=0 & \text { on } \partial \Omega\end{cases}
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with $p \in\left[1, p_{s}\right), p_{s}=\frac{d+2}{d-2},\left\|U_{p}\right\|_{p+1}=1$, so that $\left\|\nabla U_{p}\right\|_{2}^{2}=\lambda_{p}$. Then as $p \rightarrow 1$, one has $\lambda_{p} \rightarrow \lambda_{1}, U_{p} \rightarrow \Phi_{1}$ in $\mathrm{L}^{\infty}(\Omega), \nabla U_{p} \rightarrow \nabla \Phi_{1}$ in $\left(\mathrm{L}^{2}(\Omega)\right)^{d}$. Besides, there exist two explicit constants $0<c_{0}<c_{1}$ such that

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Moreover, there exists constants $0<\widetilde{k}_{0}(p) \leq \widetilde{k}_{1}(p)$ such that $\widetilde{k}_{i}(p) \rightarrow 1$ as $p \rightarrow 1^{+}$, such that

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## The End

## Thank you!!!

The kind of argument outlined here can be used to study other related problems. For example consider positive solutions to the fast diffusion equation

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\dot{u}=\Delta u^{m}, \quad \text { on } \mathbb{H}^{n}
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where

- $m \in\left(m_{s}, 1\right]$
- $\mathbb{H}^{n}$ is the hyperbolic space and $\Delta$ the corresponding Riemannian Laplacian.

Recall that, on the hyperbolic space, both the Sobolev inequality and the $\mathrm{L}^{2}$-Poincaré inequality hold, so that the $\mathrm{L}^{2}$-spectrum of $-\Delta$ is $\left[\frac{(n-1)^{2}}{4},+\infty\right)$.
Certain classes of positive solutions do vanish in a finite time (Bonforte, G. Vazquez, JEE 2008). There are some rough estimates on the extinction time there.

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It can be expected that the asymptotics of solutions is related to solutions, if any, of the elliptic problem $-\Delta u=u^{1 / m}$ (up to rescalings). No results on this till recently, but:

- Mancini-Sandeep (Annali Pisa, 2008) have shown that there exist exactly one solution $U$ to the elliptic problem. It is radial, and it has finite energy, namely it belongs to $W^{1,2}\left(\mathbb{H}^{n}\right)$. It decays at infinity as $c e^{-(n-1) r}, r$ being the Riamannian distance from the given point. There are infinitely many other radial positive solutions, but they have infinite energy. Notice that $U^{1 / m} \in \mathrm{~L}^{1}$.
- M.B., F. Gazzola, G. Grillo and J. L. Vázquez have just proved that there is no other positive radial solution apart the ones found above, and that all of them apart $U$ decay polynomially at infinity, hence they do not belong to $L^{q}$ for any $q \neq \infty$ (recall that the Riemannian measure has a density whose radial part is $\left.e^{(n-1) r}\right)$.

Hence the asymptotics of solutions to the fast diffusions should be related to the separate variable solution $\mathcal{U}(t, x)=U(x)[(T-t) / T]^{1 /(1-m)}$. Presently under investigation using the above methods, as a part of a more general study of nonlinear diffusions on manifolds.

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