# A pointwise estimate for fractionary derivatives with applications to partial differential equations

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This article emphasizes the role played by a remarkable pointwise inequality satisfied by fractionary derivatives in order to obtain maximum principles and Lp-decay of solutions of several interesting partial differential equations. In particular, there are applications to quasigeostrophic flows, in two space variables with critical viscosity, that model the Eckman pumping [see Baroud, Ch. N., Plapp, B. B., She, Z. S. & Swinney, H. L. (2002) Phys. Rev. Lett. 88, 114501 and Constantin, P. (2002) Phys. Rev. Lett. 89, 184501].

he decay in time of the spatial  $L^p$ -norm,  $1 \le p \le \infty$ , is an The decay in time of the spatial 2 morns, 1 important objective in order to understand the behavior of solutions of partial differential equations. The purpose of this article is to analyze the following pointwise inequality,  $2\theta \Lambda^{\alpha} \theta(x) \geq \Lambda^{\alpha} \theta^{2}(x)$ , valid for fractionary derivatives in  $\mathbb{R}^{n}$ ,  $n \geq 1$  $1, 0 \le \alpha \le 2$ , together with its applications to several maximum principle and decay estimates. In particular, it is applied to the quasigeostrophic equation with critical viscosity

$$\theta_t + R(\theta) \cdot \nabla^{\perp} \theta = -\kappa \Lambda \theta$$

where  $\nabla^{\perp}\theta = (-\partial\theta/\partial x_2, \partial\theta/\partial x_1)$ ,  $R(\theta) = (R_1(\theta), R_2(\theta))$ , and  $R_j$  denotes the  $j^{th}$ -Riesz transform in  $R^2$  (see refs. 1–6).

Given a weak solution  $\theta(x, t)$  (obtained as limit of solutions of the equations

$$\theta_t^{\varepsilon} + R(\theta^{\varepsilon}) \cdot \nabla^{\perp} \theta^{\varepsilon} = -\kappa \Lambda \theta^{\varepsilon} + \varepsilon \Delta \theta^{\varepsilon}$$

with the same initial data  $\theta_0$ , when the artificial viscosity  $\varepsilon$  tends to zero), it is proved that  $\|\theta(\cdot, t)\|_{L^p}$ ,  $1 \le p \le \infty$ , decays and, furthermore, there is a time  $T = T(\kappa, \theta_0) < \infty$  after which  $\theta$ becomes regular.

In this article, we describe the main ideas of the proofs, together with some of the heuristic arguments. The complete details will appear elsewhere.

### **Pointwise Estimate**

The nonlocal operator  $\Lambda = (-\Delta)^{1/2}$  is defined with the Fourier transform by  $\widehat{\Lambda}f(\xi) = |\xi|\widehat{f}$ , where  $\widehat{f}$  is the Fourier transform of f.

**Theorem 1.** Let  $0 \le \alpha \le 2, x \in \mathbb{R}^n, T^n \ (n = 1, 2, 3...)$  and  $\theta \in$  $C_0^2(\mathbb{R}^n)$ ,  $C^2(\mathbb{T}^n)$ . Then the following inequality holds:

$$2\theta \Lambda^{\alpha} \theta(x) \ge \Lambda^{\alpha} \theta^2(x)$$
. [1]

Proof: It is easy to check that the inequality is satisfied when  $\alpha = 0$  and  $\alpha = 2$ . For  $0 < \alpha < 2$  and  $n \ge 2$ , there are the formulas

$$\Lambda^{\alpha}\theta(x) = C_{\alpha}PV \int \frac{\left[\theta(x) - \theta(y)\right]}{|x - y|^{n + \alpha}} dy \quad x \in \mathbb{R}^n$$
 [2]

$$\Lambda^{\alpha}\theta(x) = \tilde{C}_{\alpha} \sum_{v \in \mathbb{Z}^n} PV \int_{\mathbb{T}^n} \frac{\left[\theta(x) - \theta(y)\right]}{\left|x - y - v\right|^{n+\alpha}} \, dy \quad x \in \mathbb{T}^n, \quad [3]$$

where  $C_{\alpha}$ ,  $\tilde{C}_{\alpha} > 0$ .

With Eq. 2 (and Eq. 3 in the periodic case) inequality Eq. 1 is obtained easily:

$$\begin{split} \theta \Lambda^{\alpha} \theta(x) &= C_{\alpha} P V \int \frac{\left[\theta(x)^2 - \theta(y) \theta(x)\right]}{\left|x - y\right|^{n + \alpha}} \, dy \\ &= \frac{1}{2} \, C_{\alpha} P V \int \frac{\left[\theta(y) - \theta(x)\right]^2}{\left|x - y\right|^{n + \alpha}} \, dy \\ &+ \frac{1}{2} \, C_{\alpha} P V \int \frac{\left[\theta^2(x) - \theta^2(y)\right]}{\left|x - y\right|^{n + \alpha}} \, dy \\ &\geq \frac{1}{2} \, \Lambda^{\alpha} \theta^2(x). \end{split}$$

The proof of the remainder case n = 1 is as follows. Given  $\psi(x_1)$  an application of the previous case, n=2, to the function  $\phi_{\delta}(x_1, x_2) = \psi(x_1)e^{-\pi\delta x_2^2}$  yields

$$\begin{aligned} 2\psi(x_1)\Lambda^{\alpha}\psi(x_1) &= \lim_{\delta \to 0} 2\phi_{\delta}(x_1, x_2)\Lambda^{\alpha}\phi_{\delta}(x_1, x_2) \\ &\geq \lim_{\delta \to 0} \Lambda^{\alpha}\phi_{\delta}^2(x_1, x_2) \\ &= \Lambda^{\alpha}\psi^2(x_1). \end{aligned}$$

Remark 1: The family of test functions  $x^p e^{-\delta x^2}$ ,  $\delta > 0$ , shows that the condition  $\alpha \leq 2$  cannot be improved.

Also, the hypothesis  $\theta \in C_0^2(\mathbb{R}^n)$ ,  $C^2(\mathbb{T}^n)$  is not necessary. Inequality Eq. 1 holds when  $\theta(x)$ ,  $\Lambda^{\alpha}\theta(x)$ ,  $\Lambda^{\alpha}\theta^{2}(x)$  are defined everywhere and are, respectively, the limits of the sequences  $\theta_m(x)$ ,  $\Lambda^{\alpha}\theta_m(x)$ ,  $\Lambda^{\alpha}\theta_m^2(x)$ , where  $\theta_m \in C_0^2(\mathbb{R}^n)$ ,  $C^2(\mathbb{T}^n)$  for each m.

## **Applications**

LP Decay. Let it be given the following scalar equation

$$(\partial_t + u \cdot \nabla) \theta = -\kappa \Lambda^{\alpha} \theta,$$

where the vector u satisfies either  $\nabla \cdot u = 0$  or  $u_i = G_i(\theta)$ , together with the appropriate hypothesis about regularity and decay at infinity, which will be specified each time, in order to allow the integration by parts needed in the proofs.

**Lemma 1.** If  $0 \le \alpha \le 2$  and  $\theta \in C_0^2(\mathbb{R}^n)(\mathbb{C}^2(\mathbb{T}^n))$ , it follows that

$$\int |\theta|^{p-2} \theta \Lambda^{\alpha} \theta dx \ge \frac{1}{p} \int |\Lambda^{\alpha/2} \theta^{p/2}|^2 dx,$$
 [4]

where  $p = 2^{j}$  and j is a positive integer.

Proof: An iterated application of inequality Eq. 1 yields:

$$\int |\theta|^{p-2} \theta \Lambda^{\alpha} \theta dx \ge \frac{1}{2} \int |\theta|^{p-2} \Lambda^{\alpha} \theta^{2} dx = \int |\theta|^{p-4} \theta^{2} \Lambda^{\alpha} \theta^{2} dx$$
$$\ge \frac{1}{4} \int |\theta|^{p-4} \Lambda^{\alpha} \theta^{4} dx \ge \frac{1}{2^{k-1}} \int |\theta|^{p-2^{k}} \Lambda^{\alpha} \theta^{2^{k}} dx,$$

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taking k = j - 1 and using Parseval's identity with the Fourier transform inequality Eq. 4 is obtained.

Remark 2: When  $p = 2^{j}$   $(j \ge 1)$  Lemma 1 implies the following improved estimate:

$$\frac{d}{dt} \|\theta\|_{L^p}^p = -\kappa p \int |\theta|^{p-2} \theta \Lambda^\alpha \theta dx$$
$$\leq -\kappa \int |\Lambda^{\alpha/2} \theta^{p/2}|^2 dx.$$

In the periodic case, this inequality yields an exponential decay of  $\|\theta\|_{L^p}$ ,  $1 \le p < \infty$ . For the nonperiodic case, Sobolev's embedding and interpolation produces

$$\frac{d}{dt} \|\theta\|_{L^p}^p \le -\kappa \left( \int \theta^{2p/2-\alpha} \, dx \right)^{2-\alpha/2}$$
$$\le -C(\|\theta\|_{L^p}^p)^{p-1+(\alpha/2)/p-1},$$

where  $C = C(\kappa, \alpha, p, \|\theta_0\|_1)$  is a positive constant. It then follows

$$\|\theta(\cdot, t)\|_{L^p}^p \le \frac{\|\theta_0\|_{L^p}^p}{(1 + \varepsilon Ct \|\theta_0\|_{L^p}^{p\varepsilon})^{1/\varepsilon}},$$

with  $\varepsilon = \alpha/2(p-1)$ .

Remark 3: The decay for other  $L^p$ ,  $1 , is obtained easily by interpolation. However, the <math>L^{\infty}$  decay needs further arguments that will be presented in the next section.

# **Viscosity Solutions of the Quasigeostrophic Equation**

A weak solution of

$$\theta_{\star} + R(\theta) \cdot \nabla^{\perp} \theta = -\kappa \Lambda \theta$$

will be called a viscosity solution with initial data  $\theta_0 \in H^s(R^2)(H^s(T^2))$ , s > 1, if it is the weak limit of a sequence of solutions, as  $\varepsilon \to 0$ , of the problems

$$\theta_t^{\varepsilon} + R(\theta^{\varepsilon}) \cdot \nabla^{\perp} \theta^{\varepsilon} = -\kappa \Lambda \theta^{\varepsilon} + \varepsilon \Delta \theta^{\varepsilon},$$
 [5]

with  $\theta^{\varepsilon}(x, 0) = \theta_0$ .

**Theorem 2.** Let  $\theta^{\varepsilon}$ ,  $\varepsilon > 0$ , be a solution of Eq. 5, then  $\theta^{\varepsilon}(\cdot, t) \in H^{s}$  for each t > 0 and satisfies

$$\|\theta^{\varepsilon}(\cdot,t)|_{L^{\infty}} \leq \frac{\|\theta_{0}\|_{L^{\infty}}}{1 + Ct \frac{\kappa \|\theta_{0}\|_{L^{\infty}}}{\|\theta_{0}\|_{L^{2}}}}$$

uniformly on  $\varepsilon > 0$  for all time  $t \ge 0$ . Furthermore, for  $s > \frac{3}{2}$ , there is a time  $T_1 = T_1(\kappa, \|\theta_0\|_{H^s})$  such that  $\|\Lambda^s \theta^{\varepsilon}(t)\|_{L^2} \le 2\|\Lambda^s \theta_0\|_{L^2}$  for  $0 \le t < T_1$ .

**Theorem 3.** Let  $\theta$  be a viscosity solution with initial data  $\theta_0 \in H^s$ ,  $s > \sqrt[3]{2}$ , of the equation  $\theta_t + R(\theta) \cdot \nabla^{\perp} \theta = -\kappa \Lambda \theta$  ( $\kappa > 0$ ). Then there exist two times  $T_1 \leq T_2$  depending only on  $\kappa$  and the initial data  $\theta_0$  so that:

- (i) If  $t \le T_1$  then  $\theta(\cdot, t) \in C^1([0, T_1); H^s)$  is a classical solution of the equation satisfying
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$$\|\theta(\cdot,t)\|_{H^s} \ll \|\theta_0\|_{H^s}.$$

(ii) If  $t \ge T_2$  then  $\theta(\cdot, t) \in C^1([T_2, \infty); H^s)$  is also a classical solution and  $\|\theta(\cdot, t)\|_{H^s}$  is monotonically decreasing in t, bounded by  $\|\theta_0\|_{H^s}$ , and satisfying

$$\int_{T_2}^{\infty} \|\theta\|_{H^s}^2 dt < \infty.$$

In particular, this implies that

$$\|\theta(\cdot,t)\|_{H^s} = O(t^{-1/2}) t \rightarrow \infty.$$

Sketch of the Proofs: For the  $L^{\infty}$ -decay there is the following heuristic argument. Assuming that  $\theta(\cdot, t)$  get its maximum value at the point  $x_t$ , depending smoothly on t, then the equation yields

$$\frac{\partial \theta}{\partial t}(x_t, t) = -\kappa \Lambda \theta(x_t, t) = -\kappa PV \int \frac{\left[\theta(x_t, t) - \theta(y, t)\right]}{|x_t - y|^3} dy.$$

And the decay is obtained because

$$PV \int \frac{[\theta(x_t, t) - \theta(y, t)]}{|x_t - y|^3} dy \ge C \frac{\theta^2(x_t, t)}{\|\theta(\cdot, t)\|_{L^2}}.$$

In the actual *Proof* the differentiability properties of Lipschitz functions are used in order to avoid the hypothesis about the existence of  $dx_t/dt$ .

The *Proof* of *Theorem 3* is based on both the  $L^{\infty}$ -decay and a bootstrap mechanism associated with the evolution of different Sobolev norms. A crucial ingredient is the fact that fR(f) belongs to Hardy's space  $\mathcal{H}^1$  for each  $L^2$ -function f and every odd singular integral R. A typical example of that mechanism is the following chain of inequalities

$$\begin{split} \frac{d}{dt} \|\theta^{\varepsilon}\|_{L^{2}}^{2} &= 2 \int \theta^{\varepsilon} R(\theta^{\varepsilon}) \cdot \nabla^{\perp} \theta^{\varepsilon} - 2\kappa \int \theta^{\varepsilon} \Lambda \theta^{\varepsilon} - 2\varepsilon \int |\Lambda \theta^{\varepsilon}|^{2} \\ &= -2\kappa \|\Lambda^{1/2} \theta^{\varepsilon}\|_{L^{2}}^{2} - 2\varepsilon \|\Lambda \theta^{\varepsilon}\|_{L^{2}}^{2} \leq -2\kappa \|\Lambda^{1/2} \theta^{\varepsilon}\|_{L^{2}}^{2} \\ \frac{d}{dt} \|\Lambda^{1/2} \theta^{\varepsilon}\|_{L^{2}}^{2} &= 2 \int \Lambda^{1/2} \theta^{\varepsilon} \Lambda^{1/2} (R(\theta^{\varepsilon}) \cdot \nabla^{\perp} \theta^{\varepsilon}) - 2\kappa \|\Lambda \theta^{\varepsilon}\|_{L^{2}}^{2} \\ &- 2\varepsilon \|\Lambda^{3/2} \theta^{\varepsilon}\|_{L^{2}}^{2} \\ &\leq (C \|\theta^{\varepsilon}(\cdot, t)\|_{L^{\infty}} - 2\kappa) \|\Lambda \theta^{\varepsilon}\|_{L^{2}}^{2} \\ \frac{d}{dt} \|\Lambda \theta^{\varepsilon}\|_{L^{2}}^{2} &\leq (C \|\Lambda \theta^{\varepsilon}\|_{L^{2}} - \kappa) \|\Lambda^{3/2} \theta^{\varepsilon}\|_{L^{2}}^{2} \\ \frac{d}{dt} \|\Lambda^{3/2} \theta^{\varepsilon}\|_{L^{2}}^{2} &\leq (C \|\theta^{\varepsilon}\|_{L^{\infty}} - \kappa) \|\Delta \theta^{\varepsilon}\|_{L^{2}}^{2} \end{split}$$

valid for some universal constant C, uniformly with respect to the artificial viscosity  $\varepsilon$ .

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